

Complex-scaling method for the plasmonic resonances of particles with corners

DeFI working group (INRIA-CMAP)

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in collaboration with

Anne-Sophie Bonnet-Ben Dhia¹ and Christophe Hazard¹

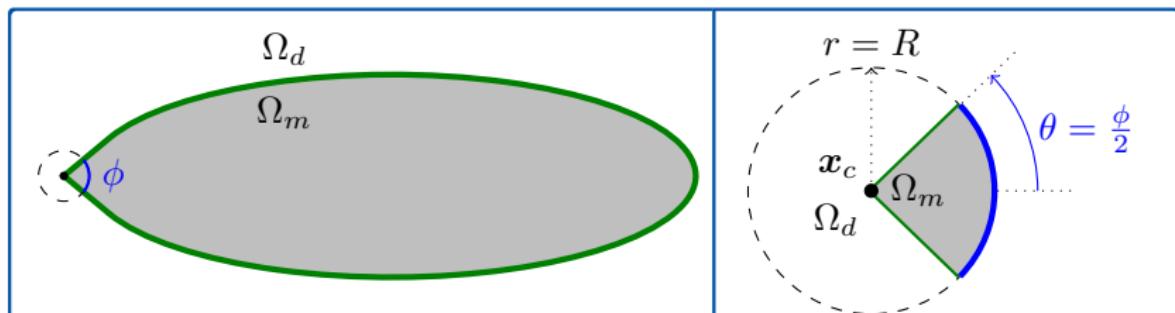
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Objective

Context: Waves in “negative” materials, i.e. $\epsilon_m(\omega) < 0$, $\mu_m(\omega) < 0$

Particle: Boundary $\partial\Omega_m$ has one straight corner of angle $\phi \in (0, 2\pi) \setminus \pi$



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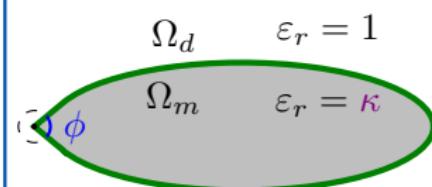
Plasmonic Eigenvalue Problem (PEP)

Find $(u, \kappa) \in U \times \mathbb{C}$ such that

$$\operatorname{div} [\varepsilon_r^{-1} \nabla u] = 0 \quad \text{in } \mathcal{D}'(\mathbb{R}^2)$$

with piecewise-constant permittivity:

$$\varepsilon_r = \kappa \mathbf{1}_{\Omega_m} + \mathbf{1}_{\Omega_d}$$



⚠ Spectral parameter: “contrast” $\kappa := \frac{\epsilon_m(\omega)}{\epsilon_d}$.

PEP is self-adjoint in $L^2(\mathbb{R}^2)$ with $\kappa < 0 \Rightarrow$ **sign-changing interface**.

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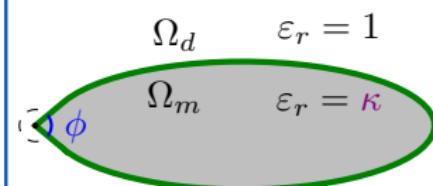
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Objective Investigate existence of solutions $(u_{\text{res}}, \kappa_{\text{res}})$ with

$$\kappa_{\text{res}} \in \mathbb{C} \setminus \mathbb{R}, \quad u_{\text{res}} \notin L^2_{\text{loc}}(\mathbb{R}^2).$$

Next: why a corner?

Smooth boundary: H^1 plasmons (discrete spectrum)

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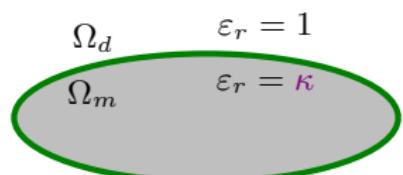
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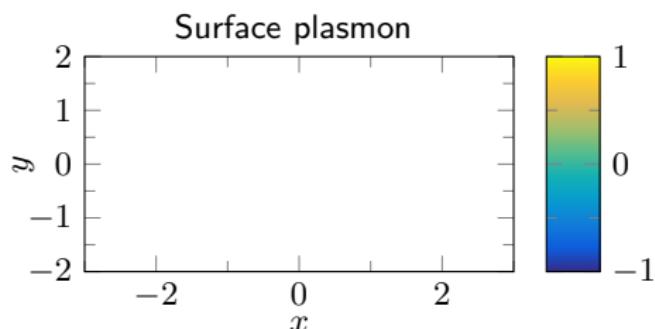
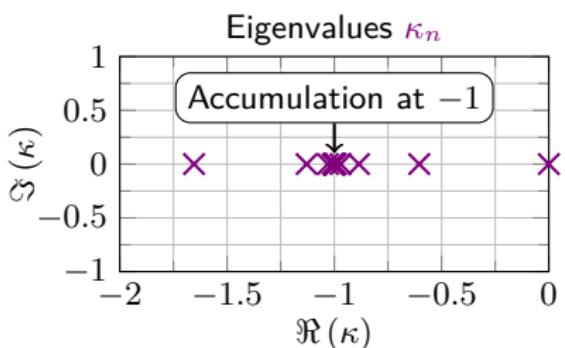
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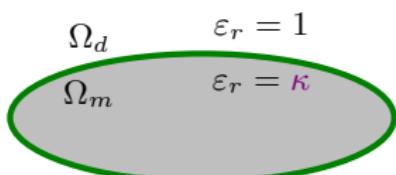
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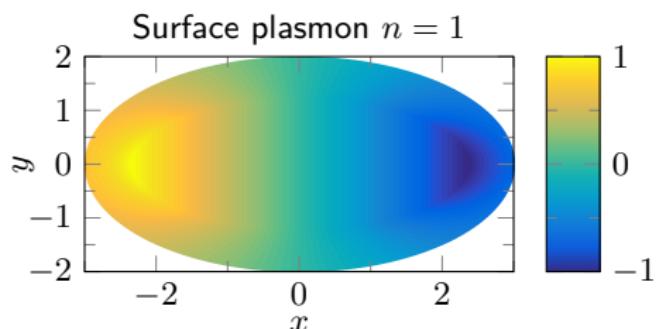
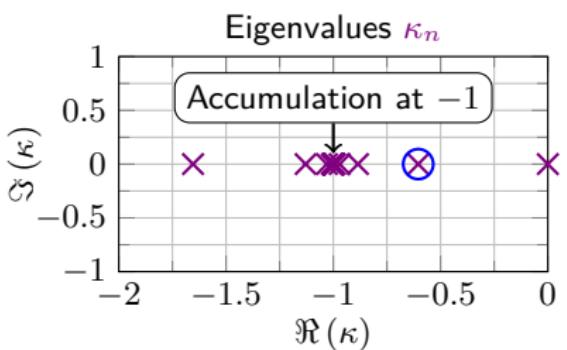
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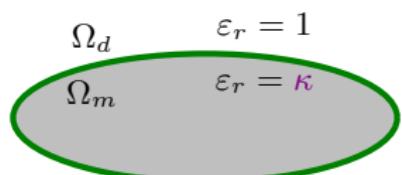
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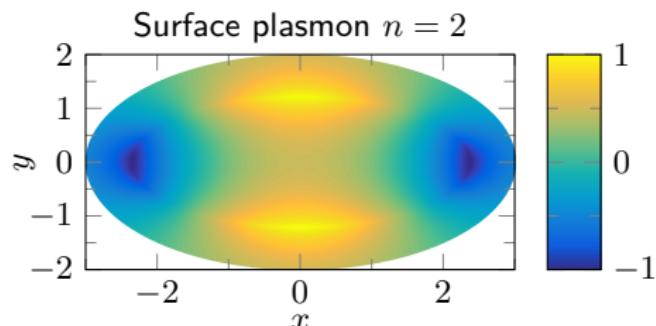
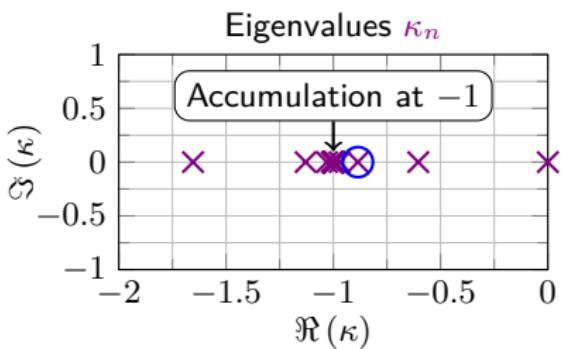
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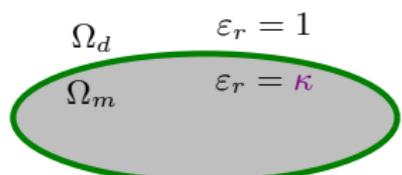
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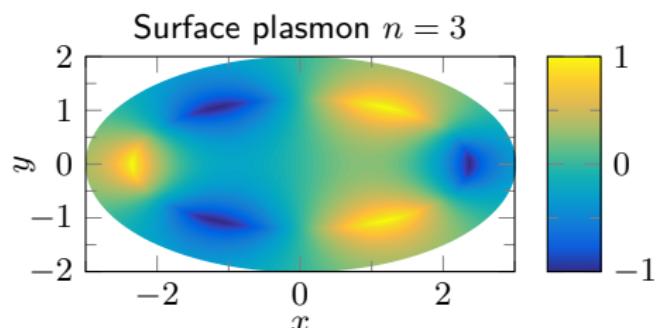
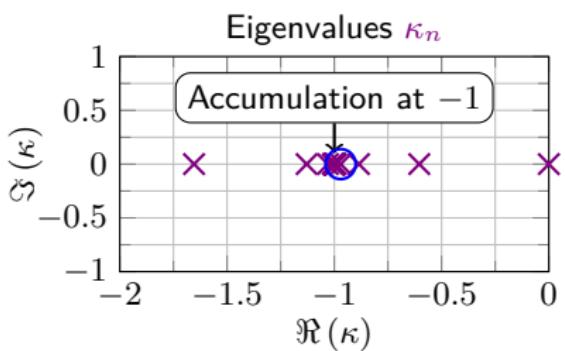
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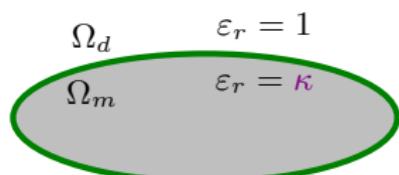
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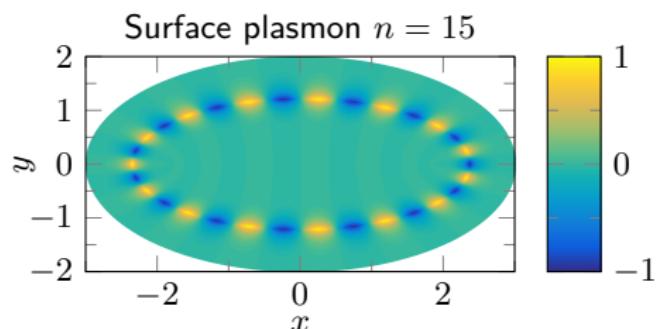
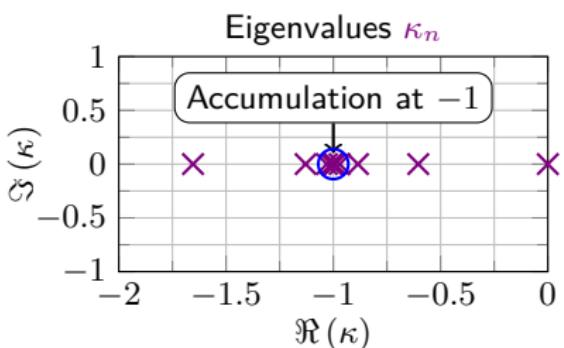
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Next: effect of corner on plasmons?

Corner: strongly-oscillating plasmons (essential spectrum)

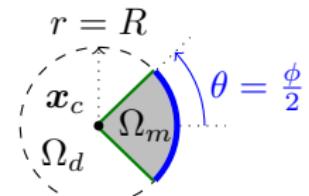
Let $\partial\Omega_m$ have one corner at x_c and $D := \{|\mathbf{x} - \mathbf{x}_c| < R\}$.

Even (odd) local solutions with separated variables are:

$$u_{\eta}^{e(o)}(r, \theta) := r^{i\eta} \times \Phi_{\eta}^{e(o)}(\theta)$$

where

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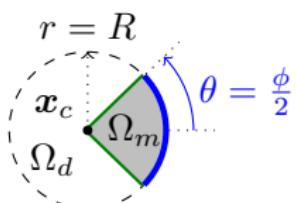
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Analysis of $f_{\phi}^{e(o)}$ yields **critical interval** $I_c = I_c^e \cup I_c^o$: (Bonnet-Ben Dhia et al. 2016)

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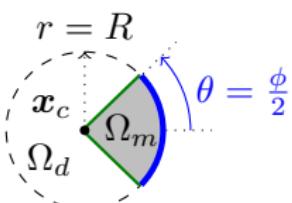
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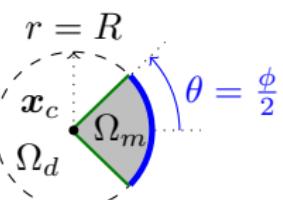
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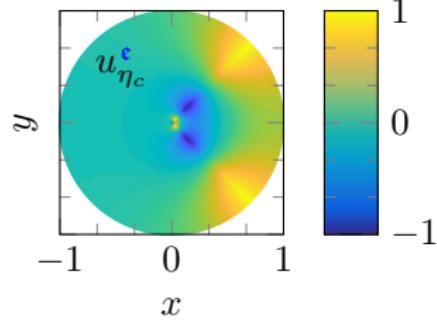
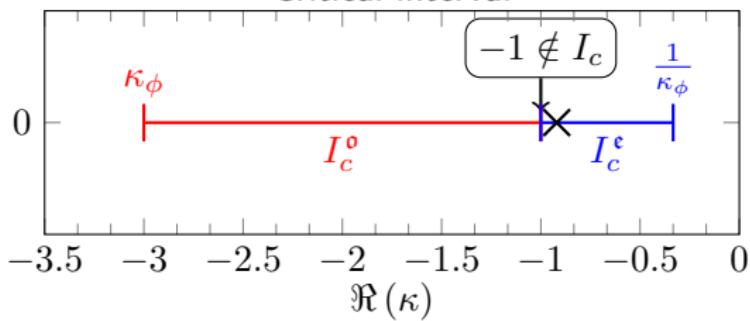
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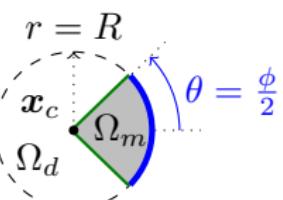
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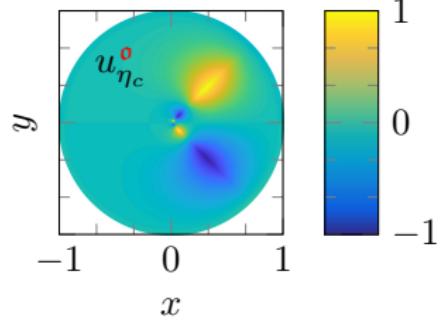
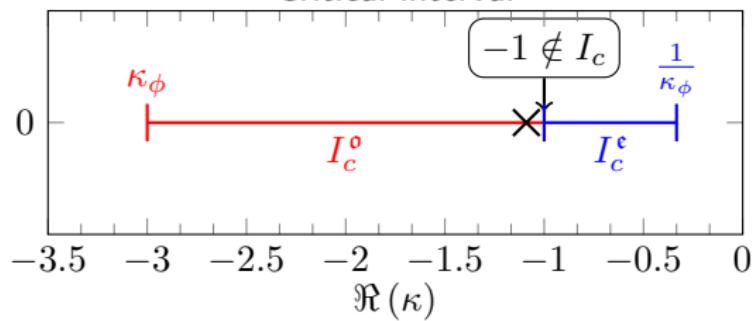
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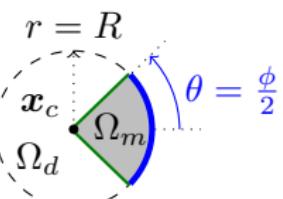
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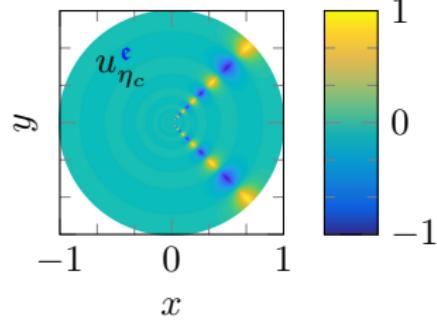
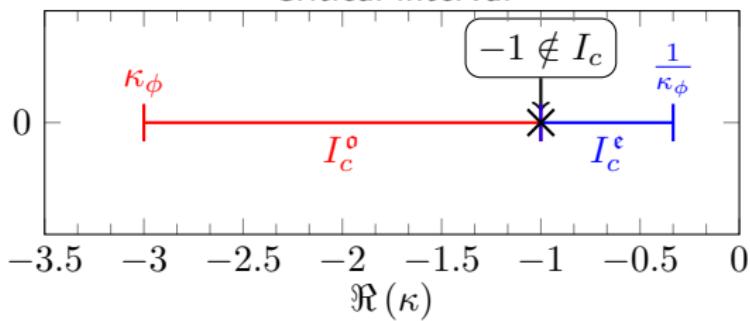
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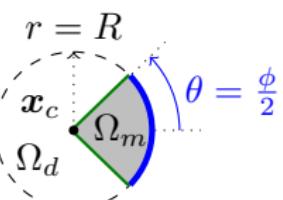
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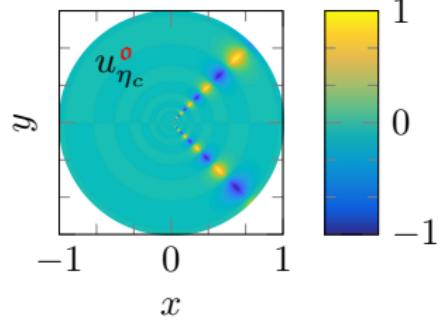
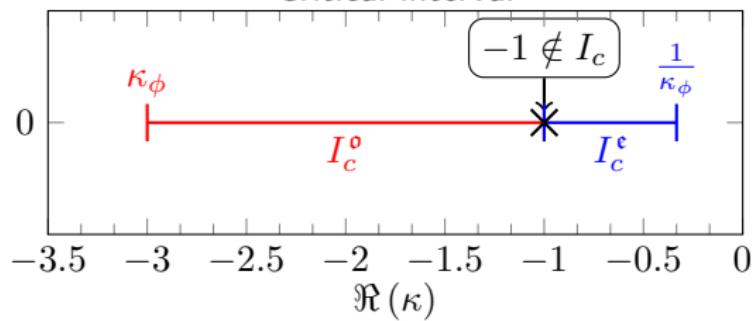
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Definition of complex plasmonic (CP) resonances: strategy

Tentative definition: A contrast $\kappa \in \mathbb{C} \setminus \mathbb{R}$ is a *CP resonance* if there is $u_{\text{res}} : \mathbb{R}^2 \rightarrow \mathbb{C}$ such that (u_{res}, κ) solves the PEP.

⚠ $u_{\text{res}} \notin L^2_{\text{loc}}(\mathbb{R}^2)$, since the PEP is self-adjoint.

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$$\widehat{H}_\phi^{\text{e}(\text{o})}(\kappa) := \{\eta \in H_\phi^{\text{e}(\text{o})}(\kappa) \mid \Im(\eta) < 0\}, \quad \widehat{H}_\phi(\kappa) := \widehat{H}_\phi^{\text{e}}(\kappa) \cup \widehat{H}_\phi^{\text{o}}(\kappa).$$

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Starting point: Let $\kappa \in \mathbb{C}^+$ and $u(\kappa) \in H^1(D)$ local solution: $\forall \eta_* < 0$,

$$u(r, \theta; \kappa) \underset{r \rightarrow 0}{=} c_0 + \sum_{\mathfrak{p} \in \{\textcolor{blue}{e}, \textcolor{red}{o}\}} \sum_{\substack{\eta \in \widehat{H}_\phi^{\mathfrak{p}}(\kappa) \\ \eta_* < \Im(\eta)}} c_\eta^\mathfrak{p} u_\eta^\mathfrak{p}(r, \theta) + \mathcal{O}(r^{-\eta_*}) \quad (1)$$

Strategy Obtain a continuation of (1) by studying the continuation to \mathbb{C}^- of the map

$$\mathbb{C}^+ \ni \kappa \mapsto \widehat{H}_\phi(\kappa).$$

Continuation of set of stable zeros: key results

Zeros: $\mathbb{C} \ni \kappa \mapsto H_\phi^{\textcolor{blue}{c}(\textcolor{red}{o})}(\kappa)$, Stable zeros: $\mathbb{C}^+ \ni \kappa \mapsto \widehat{H}_\phi^{\textcolor{blue}{c}(\textcolor{red}{o})}(\kappa)$

Lemma (analyticity). For any $\kappa \in U := \mathbb{C} \setminus \{\kappa_\phi, 1/\kappa_\phi\}$, each element η of $H_\phi^{\textcolor{blue}{c}(\textcolor{red}{o})}(\kappa)$ depends analytically upon $\kappa \in U$.

The map $\mathbb{C} \ni \kappa \mapsto H_\phi^{\textcolor{blue}{c}(\textcolor{red}{o})}(\kappa)$ is single-valued, but...

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Proposition. The map $\mathbb{C}^+ \ni \kappa \mapsto \widehat{H}_\phi^{e(o)}(\kappa)$ has three branch points:

$\kappa = \kappa_\phi$ (algebraic), $\kappa = -1$ (logarithmic), $\kappa = 1/\kappa_\phi$ (algebraic).

Proof. Follows from asymptotic expansions of the critical exponent η_c . □

By crossing \mathbb{R} **once** we obtain **three continuations** of $\widehat{H}_\phi(\kappa)$ to \mathbb{C}^- :

$$\widehat{H}_\phi(\kappa), \widehat{H}_\phi^{|e}(\kappa), \widehat{H}_\phi^{|o}(\kappa).$$

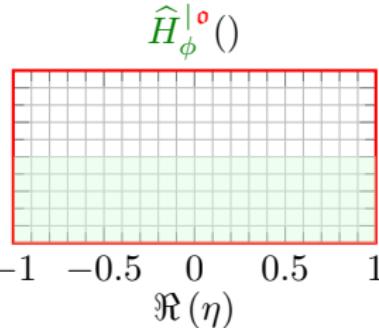
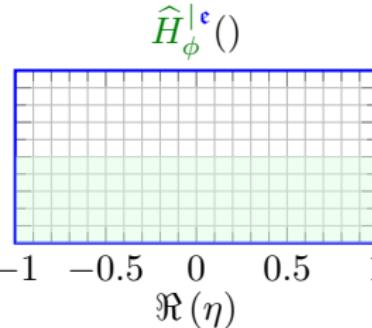
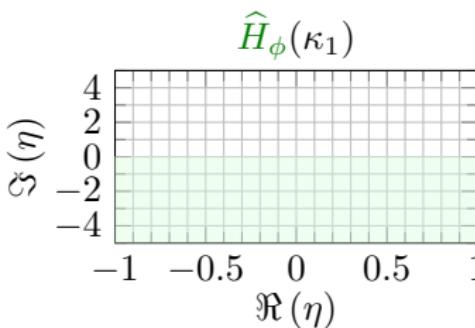
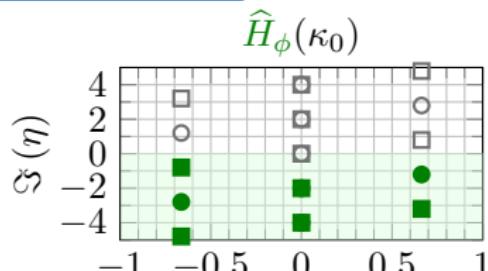
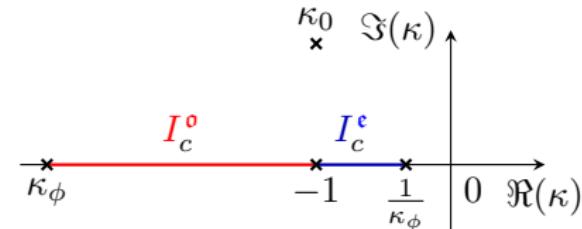
Next: let us plot each of these three continuations.

Continuation of set of stable zeros: illustration

Let's compare three paths satisfying

$$\Gamma : (0, 1) \rightarrow \mathbb{C}, \quad \Gamma(0) = \kappa_0, \quad \Gamma(1) = \kappa_1$$

We track $\hat{H}_\phi(\kappa)$ (\bullet , \blacksquare) as κ moves.



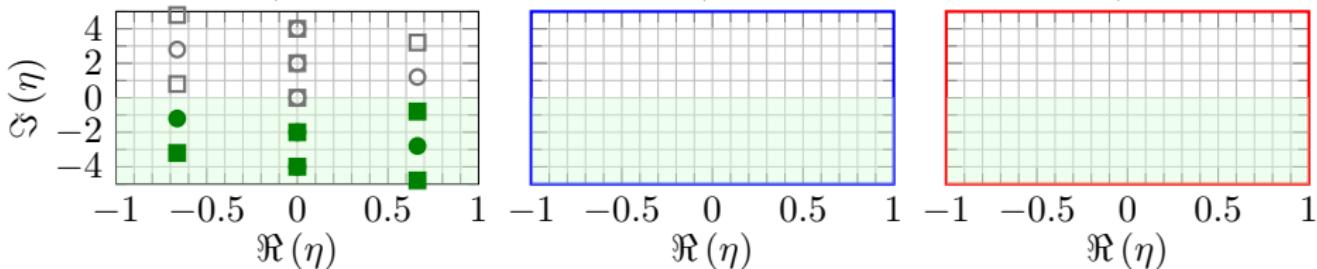
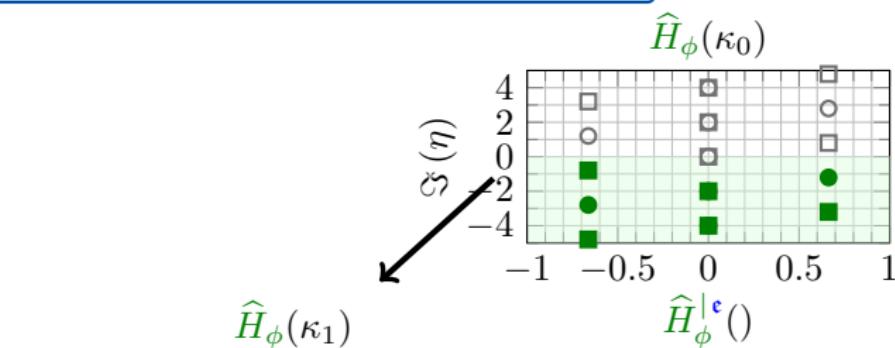
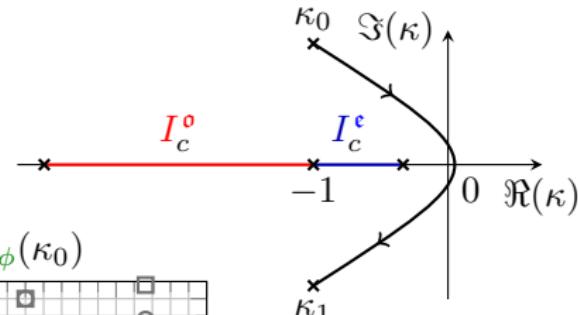
Next: definition of complex plasmonic resonances using $\hat{H}_\phi^{|e(o)}$.

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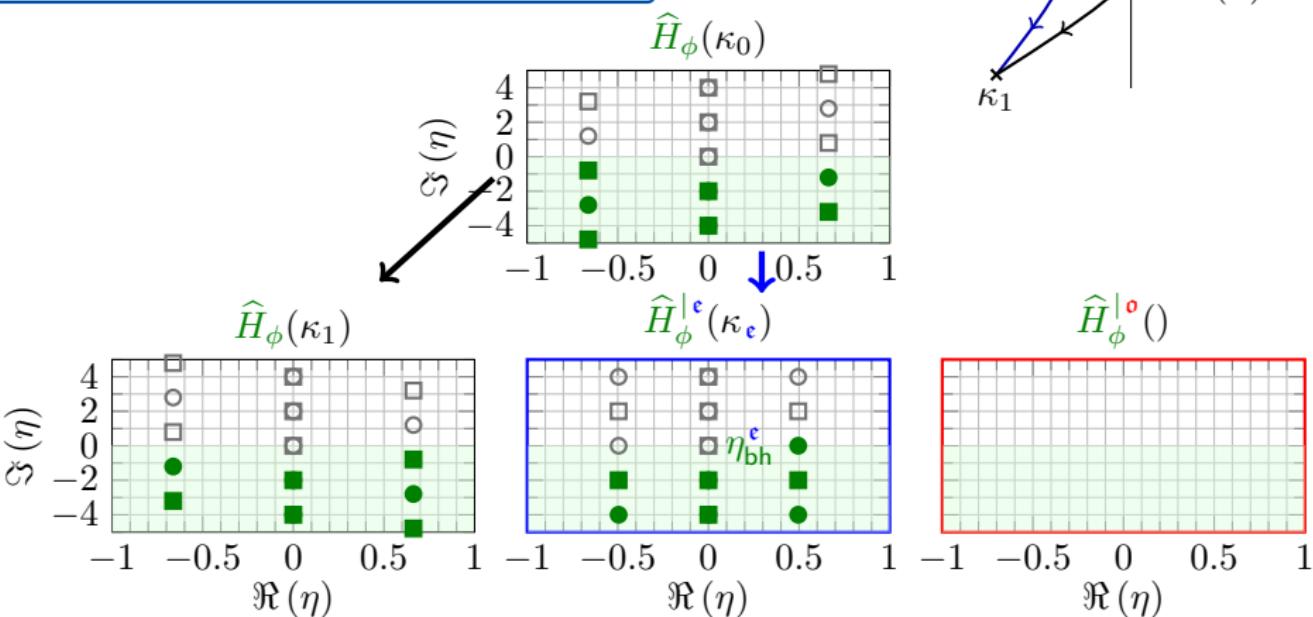
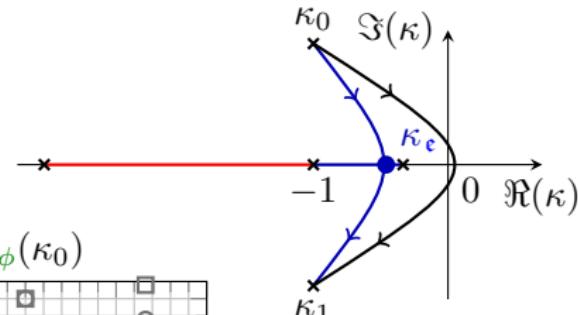
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Continuation of set of stable zeros: illustration

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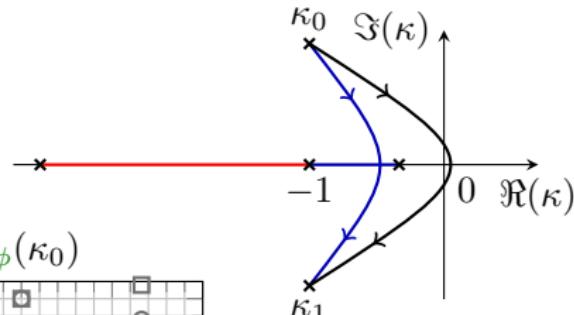
Next: definition of complex plasmonic resonances using $\hat{H}_\phi^{(c)}$.

Continuation of set of stable zeros: illustration

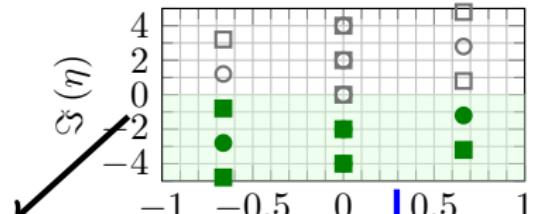
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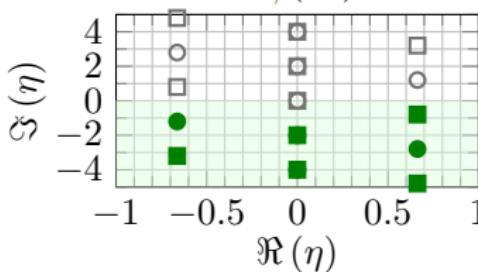
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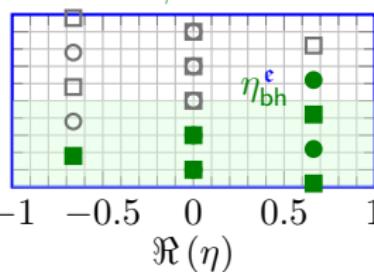
$\hat{H}_\phi(\kappa_0)$



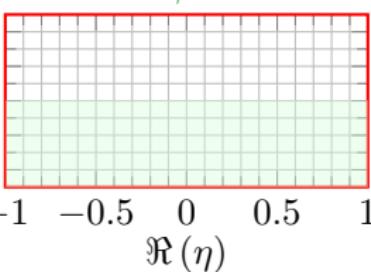
$\hat{H}_\phi(\kappa_1)$



$\hat{H}_\phi^{\mid c}(\kappa_1)$



$\hat{H}_\phi^{\mid o}()$



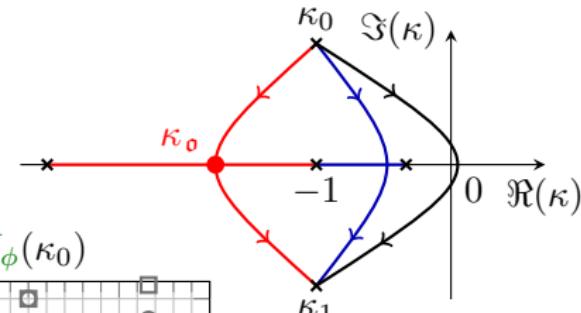
Next: definition of complex plasmonic resonances using $\hat{H}_\phi^{\mid c(o)}$.

Continuation of set of stable zeros: illustration

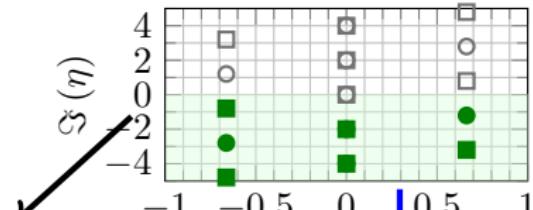
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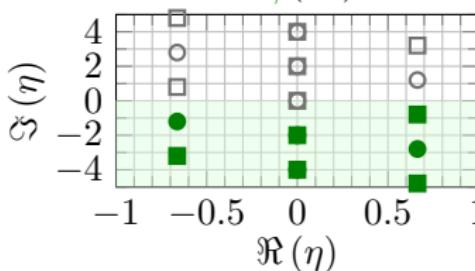
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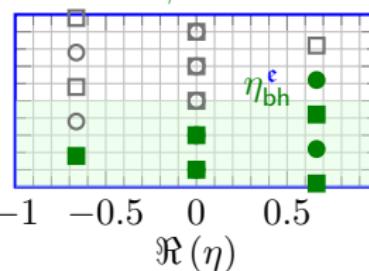
$\hat{H}_\phi(\kappa_0)$



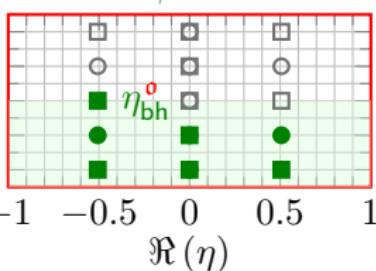
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$\hat{H}_\phi^{\mid o}(\kappa_o)$



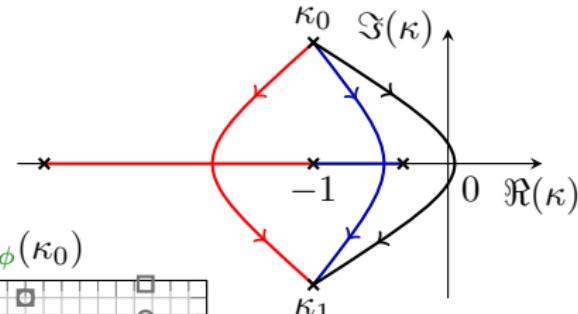
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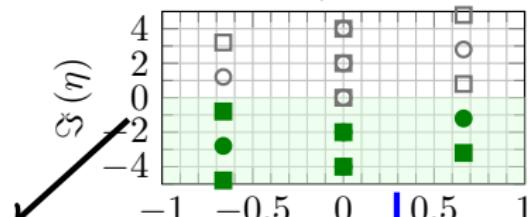
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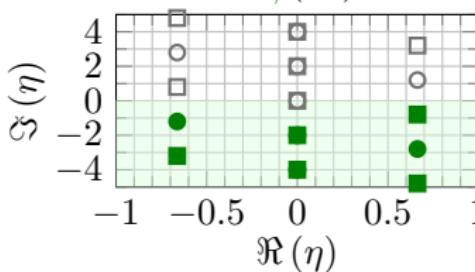
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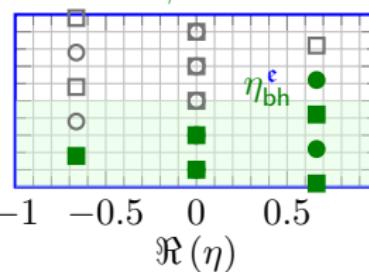
$\hat{H}_\phi(\kappa_0)$



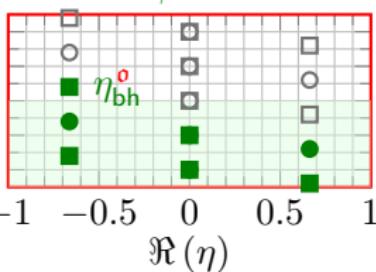
$\hat{H}_\phi(\kappa_1)$



$\hat{H}_\phi^{\mid c}(\kappa_1)$



$\hat{H}_\phi^{\mid o}(\kappa_1)$



Next: definition of complex plasmonic resonances using $\hat{H}_\phi^{\mid c(o)}$.

Definition of complex plasmonic (CP) resonances

The continuation of $\widehat{H}_\phi(\kappa)$ across I_c^e is

$$\widehat{H}_\phi^{|\text{e}}(\kappa) = \widehat{H}_\phi^{\text{e}|\text{e}}(\kappa) \cup \widehat{H}_\phi^{\text{o}|\text{e}}(\kappa).$$

This suggests...

Even-critical complex plasmonic resonance:

Contrast $\kappa_{\text{res}} \in \mathbb{C}^-$ associated with $u_{\text{res}} \notin L^2_{\text{loc}}(\mathbb{R}^2)$.

$$u_{\text{res}}(r, \theta) \underset{r \rightarrow 0}{=} c_0 + \sum_{\mathfrak{p} \in \{\text{e}, \text{o}\}} \sum_{\begin{subarray}{c} \eta \in \widehat{H}_\phi^{\mathfrak{p}|\text{e}}(\kappa_{\text{res}}), \\ \eta_\star < \Im(\eta) \end{subarray}} c_\eta^{\mathfrak{p}} u_\eta^{\mathfrak{p}}(r, \theta) + \mathcal{O}(r^{-\eta_\star})$$

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$$\underset{r \rightarrow 0}{\sim} c_0 + c_\eta^{\text{e}} r^{i\eta} \Phi_\eta^{\text{e}}(\theta) \quad \text{with } \eta = \eta_{\text{bh}}^{\text{e}}(\kappa_{\text{res}}),$$

where the blow-up rate is

$$\Im(\eta_{\text{bh}}^{\text{e}}(\kappa_{\text{res}})) > 0.$$

Next: applicability of corner complex scaling (Bonnet-Ben Dhia, Carvalho, Chesnel, and Ciarlet 2016).

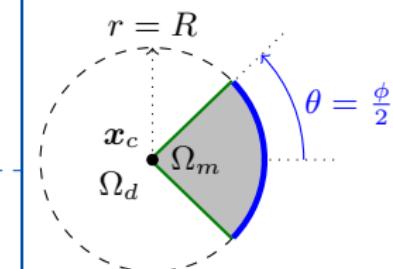
Corner complex scaling: formulation (Bonnet-Ben Dhia et al. 2016)

(PEP) Find $(u, \kappa) \in U \times \mathbb{C}$:

$$\frac{1}{r^2} r \partial_r [\varepsilon^{-1} r \partial_r u] + \frac{1}{r^2} \partial_\theta [\varepsilon^{-1} \partial_\theta u] = 0.$$

Local solutions

$$u_{\eta}^{\epsilon(\sigma)}(r, \theta) = r^{i\eta} \Phi_{\eta}^{\epsilon(\sigma)}(\theta), \quad \eta \in H_{\phi}^{\epsilon(\sigma)}(\kappa)$$



Principle: Let $\alpha \in \mathbb{C}$. Define a non self-adjoint “PEP α ” such that:

κ complex plasmonic resonance of PEP \iff κ eigenvalue of PEP α .

Intuition: Pick α such that if $u_{\text{res}} \sim r^{i\eta}$ then $u_{\text{res}, \alpha} \sim r^{i\frac{\eta}{\alpha}}$ with $\Im\left(\frac{\eta}{\alpha}\right) < 0$.

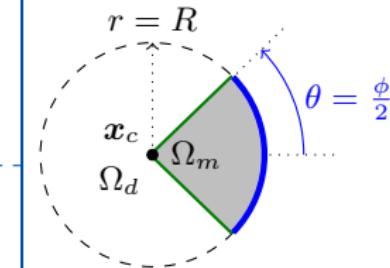
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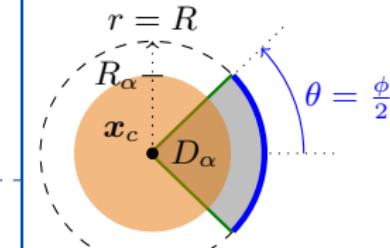
(PEP- α) Find $(u, \kappa) \in U \times \mathbb{C}$:

$$\frac{\alpha(r)}{r^2} r \partial_r [\varepsilon^{-1} \alpha(r) r \partial_r u] + \frac{1}{r^2} \partial_\theta [\varepsilon^{-1} \partial_\theta u] = 0$$

with purely-radial scaling $\alpha(r) = \alpha \mathbb{1}_{D_\alpha} + \mathbb{1}_{\mathbb{R}^2 \setminus \overline{D_\alpha}}$.

Local solutions (same dispersion relation)

$$u_{\eta, \alpha}^{\epsilon(\sigma)}(r, \theta) = r^{i\frac{\eta}{\alpha}} \Phi_{\eta}^{\epsilon(\sigma)}(\theta), \quad \eta \in H_{\phi}^{\epsilon(\sigma)}(\kappa)$$



Corner complex scaling: analysis of uncovered region

Where can complex plasmonic resonances be computed?

For any scaling $\alpha \in \mathbb{C}^*$, the **uncovered region** is

$$U_{\phi}^{\epsilon(\alpha), \alpha} := \left\{ \kappa \in \overline{\mathbb{C}^-} \mid \forall \eta \in \widehat{H}_{\phi}^{|\epsilon(\alpha)}(\kappa), \Im\left(\frac{\eta}{\alpha}\right) < 0 \right\}.$$

Plotting the uncovered region is crucial to post-process results because of...

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Plotting the uncovered region is crucial to post-process results because of...

Proposition. Let $\alpha \in \mathbb{C} \setminus \mathbb{R}$ and (u_{α}, κ) be a solution of PEP α . If

$$\kappa \in U_{\phi}^{\epsilon, \alpha}$$

then κ is an even-critical CP resonance associated with

$$u_{\text{res}}(x) := u_{\alpha}(x) \quad (x \in \mathbb{R}^2 \setminus \overline{D_{\alpha}})$$

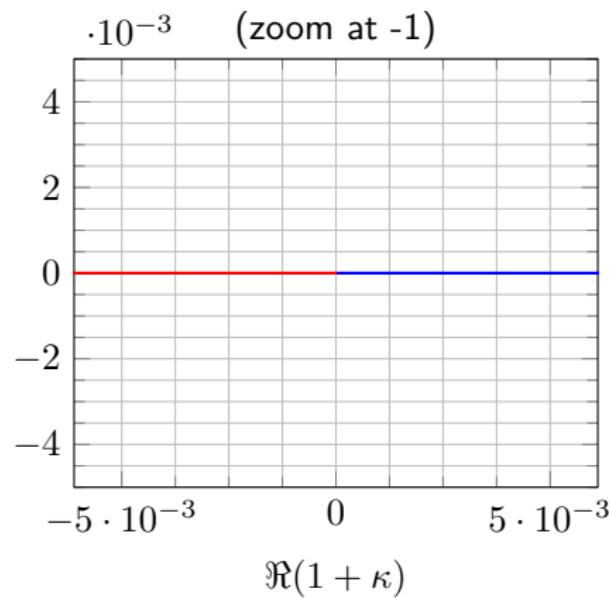
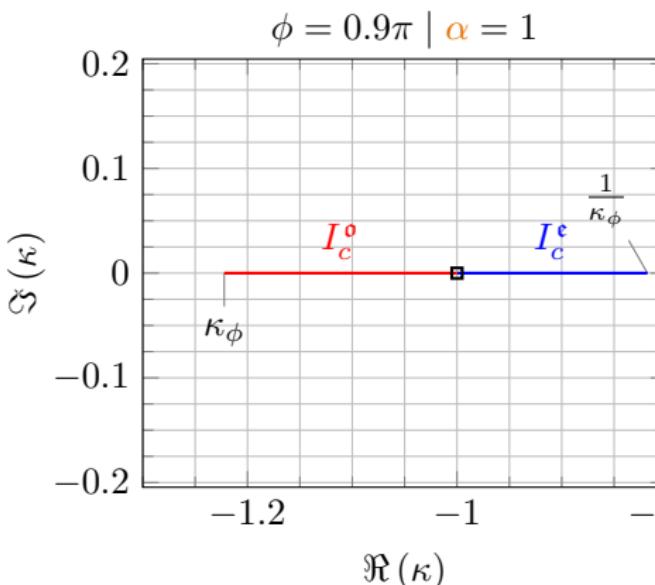
$$u_{\text{res}}(r, \theta) := u_{\alpha}(r^{\alpha}, \theta) \quad ((r, \theta) \in D_{\alpha}).$$

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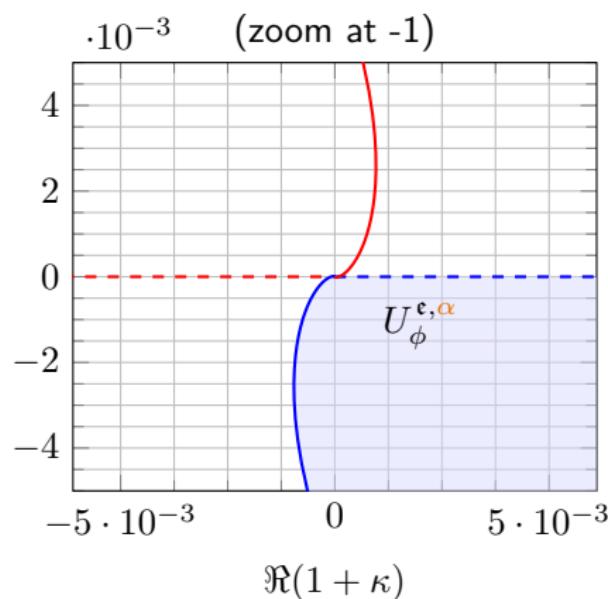
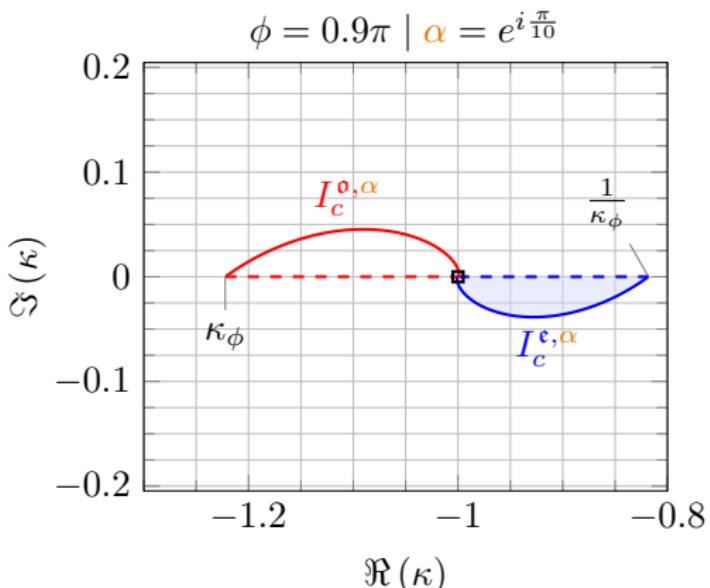
Next: discretization

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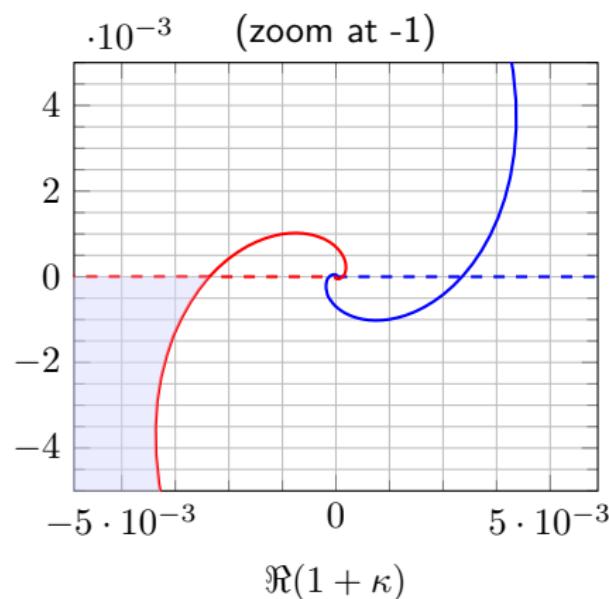
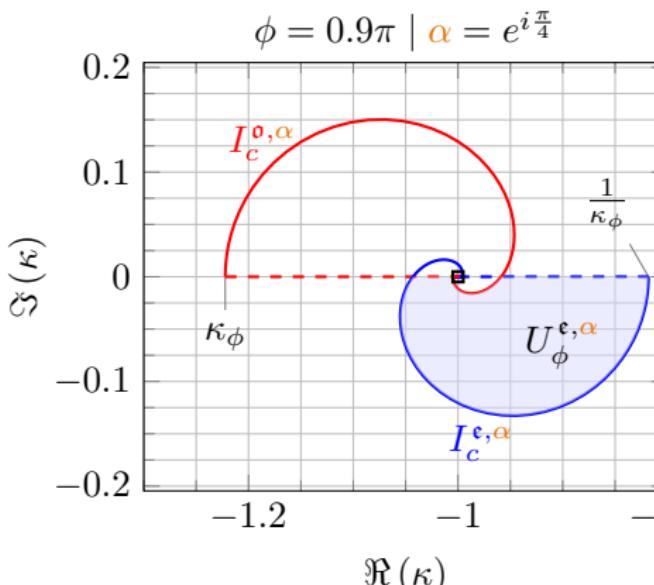
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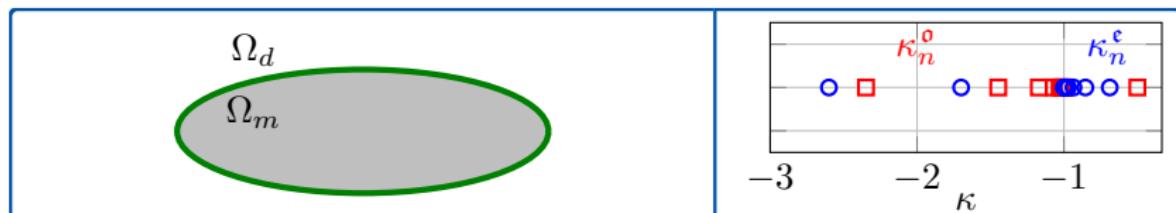
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Next: discretization

How to obtain complex plasmonic (CP) resonances?

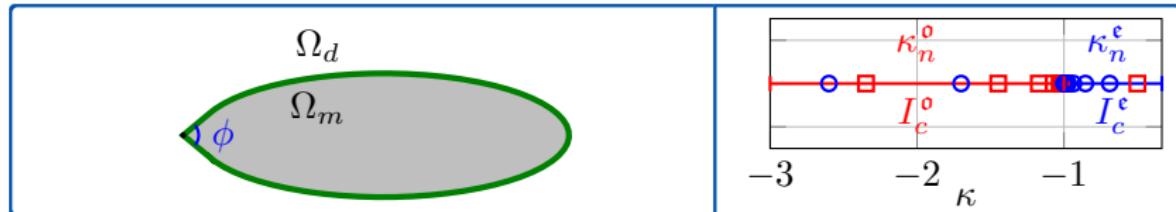
Strategy based on (Li and Shipman 2019, § 5.2): small perturbation of smooth particle Ω_m by a corner of angle ϕ



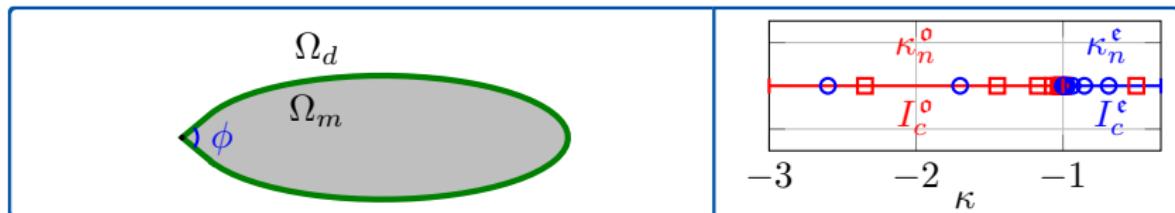
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How to obtain complex plasmonic (CP) resonances?



Perturbation mechanism described in (Li and Shipman 2019, § 5.2):

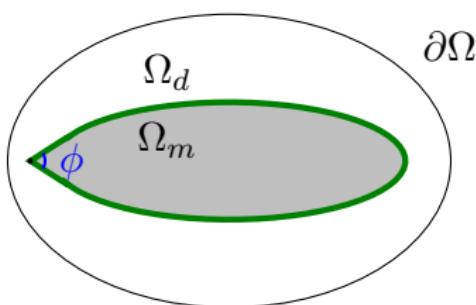
- odd eigenvalues κ_n^o in I_c^e are perturbed into embedded eigenvalues
- even eigenvalues κ_n^e in I_c^e are perturbed into complex resonances

Known results from spectral studies of the Neumann-Poincaré operator (double-layer potential) (Khavinson, Putinar, and Shapiro 2007):

- Numerical evidence of embedded eigenvalues κ_n (Helsing, Kang, and Lim 2017)
- Existence theorem for embedded eigenvalues (Li and Shipman 2019) (Perfekt 2019)

Weak formulation: without scaling

Particle: piecewise-smooth boundary $\partial\Omega_m$ of angle $\phi \in (0, \pi)$.



Weak Formulation Find $(u, \kappa) \in H_0^1(\Omega) \times \mathbb{C}$ such that

$$\forall v \in H_0^1(\Omega), \int_{\Omega_m} \nabla u(\mathbf{x}) \cdot \nabla v(\mathbf{x}) \, d\mathbf{x} = -\kappa \int_{\Omega_d} \nabla u(\mathbf{x}) \cdot \nabla v(\mathbf{x}) \, d\mathbf{x}.$$

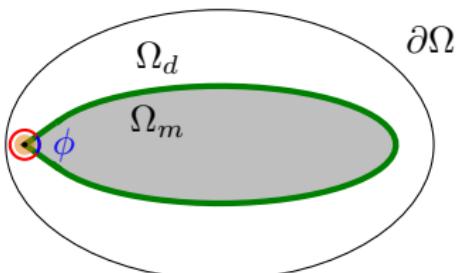
Discretization N_h -dimensional generalized eigenvalue problem

$$A_{\Omega_m} U = -\kappa A_{\Omega_d} U,$$

where each matrix is real symmetric and positive (but not definite).

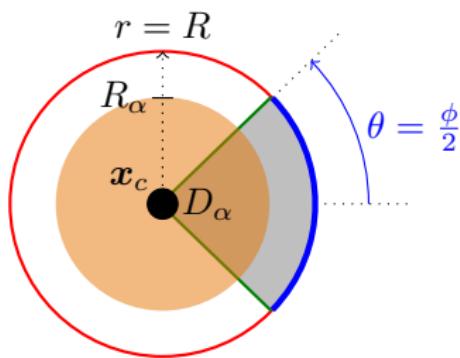
Weak formulation: with complex scaling

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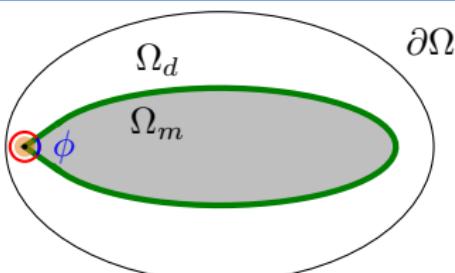
Unknown u in

$$H_e := H_0^1(\Omega)$$



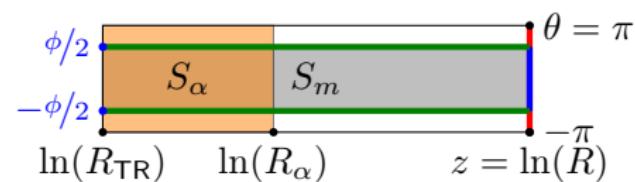
Weak formulation: with complex scaling

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Unknown u in

$$H_e := \{u \in H^1(\Omega \setminus \overline{D}) \mid u|_{\partial\Omega} = 0\}$$



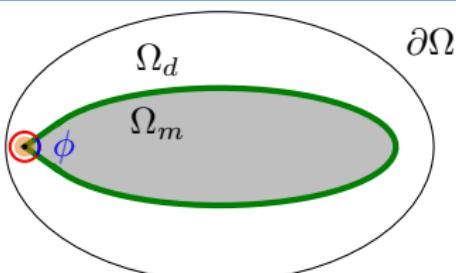
Euler coordinates ($z = \ln(r), \theta$).

Unknown \check{u} in

$$H_c := \{\check{u} \in H^1(S) \mid \check{u}(\cdot, \pi) = \check{u}(\cdot, -\pi)\}$$

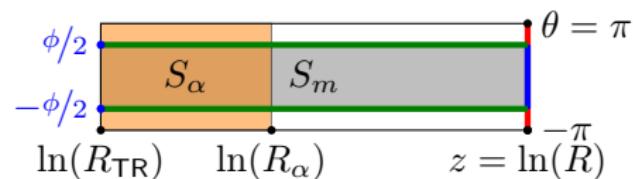
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Solution space with matching condition

$$\{(u, \check{u}) \in H_e \times H_c \mid u(x_c + R \cos \theta, R \sin \theta) = \check{u}(\ln R, \theta) \quad (\theta \in (-\pi, \pi])\}$$

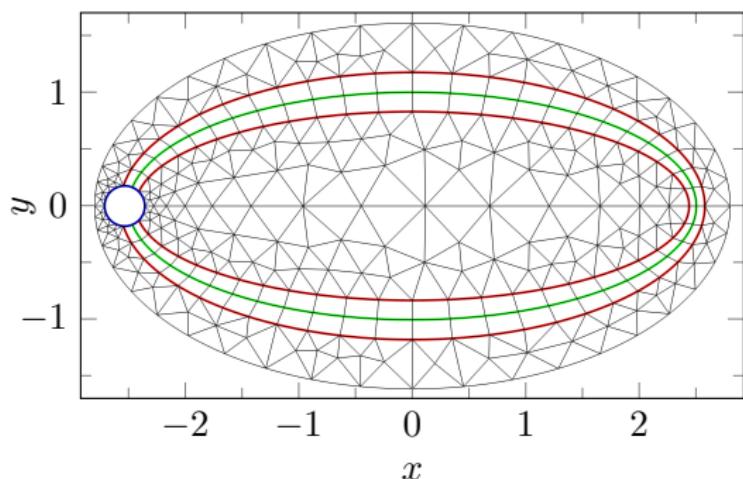
Discretization yields generalized eigenvalue problem

$$\begin{aligned} & \left[A_{\Omega_m \setminus D}^{(x,y)} + A_{S_m \setminus S_\alpha}^{(z,\theta)} + \alpha A_{S_m \cap S_\alpha}^{(z)} + \frac{1}{\alpha} A_{S_m \cap S_\alpha}^{(\theta)} \right] U = \\ & -\kappa \left[A_{\Omega_d \setminus D}^{(x,y)} + A_{S_d \setminus S_\alpha}^{(z,\theta)} + \alpha A_{S_d \cap S_\alpha}^{(z)} + \frac{1}{\alpha} A_{S_d \cap S_\alpha}^{(\theta)} \right] U \end{aligned}$$

Geometry and mesh

⚠ Local mesh symmetry at $\partial\Omega_m$ (Bonnet-Ben Dhia, Carvalho, and Ciarlet 2018).

Truncated domain $\Omega \setminus D$



$\Omega \setminus D$ (zoom)

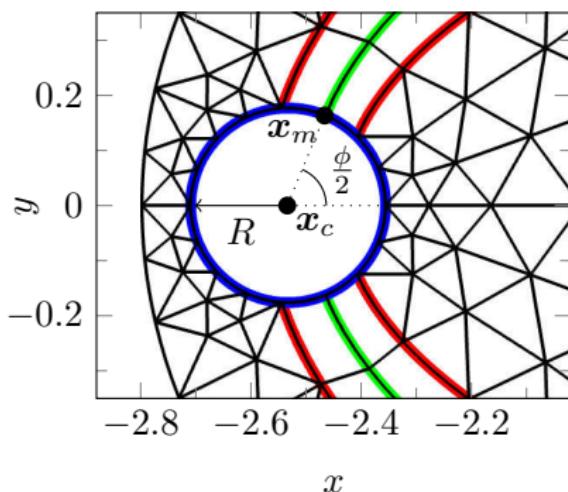


Fig. Mesh topology.



Corner-ellipse junction at x_m is C^1 :

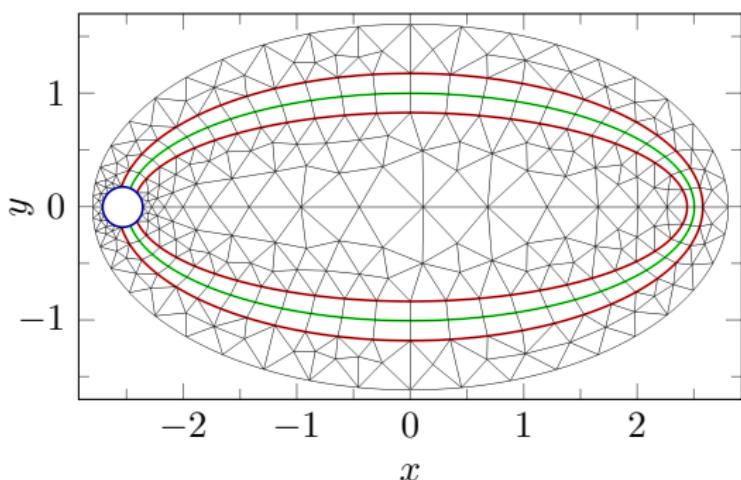
\Rightarrow boundary $\partial\Omega_m$ uniquely defined by (a_m, b_m, ϕ) .

\Rightarrow perturbation size $R := |x_c - x_m| = f(\phi, a_m, b_m)$

Geometry and mesh

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Truncated domain $\Omega \setminus D$



Corner domain S

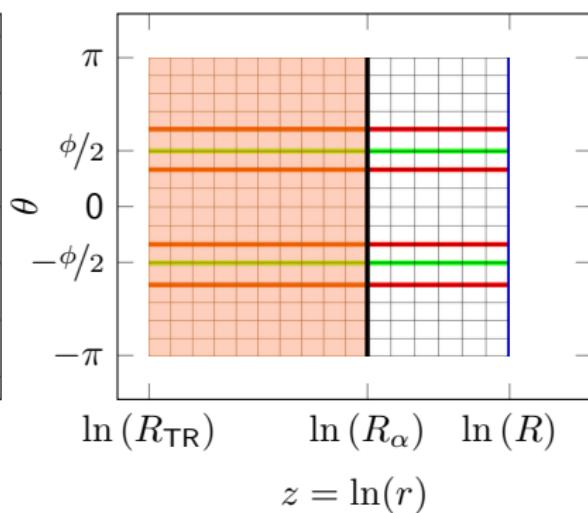


Fig. Mesh topology.



Corner-ellipse junction at x_m is C^1 :

\Rightarrow boundary $\partial\Omega_m$ uniquely defined by (a_m, b_m, ϕ) .

\Rightarrow perturbation size $R := |x_c - x_m| = f(\phi, a_m, b_m)$

Results: spectrum

$$\alpha = e^{i\frac{\pi}{40}} \mid \phi = 0.75\pi$$

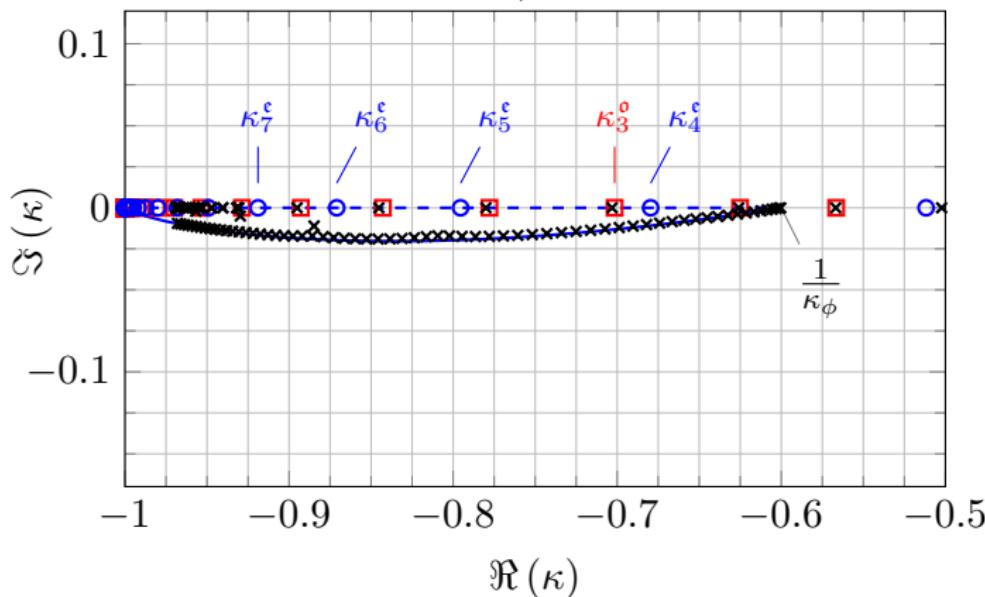


Fig. Spectrum for increasing values of $\arg(\alpha)$.

(\times): computed ($N_h = 26345$, $R_\alpha = R/2$, $R_{\text{TR}} = 10^{-50} \cdot R$),

(○): unperturbed even eigenvalues κ_n^e , (□): κ_n^o ,

(—): even critical curve $I_c^{\epsilon, \alpha}$, (—): $I_c^{\sigma, \alpha}$.

Results: spectrum

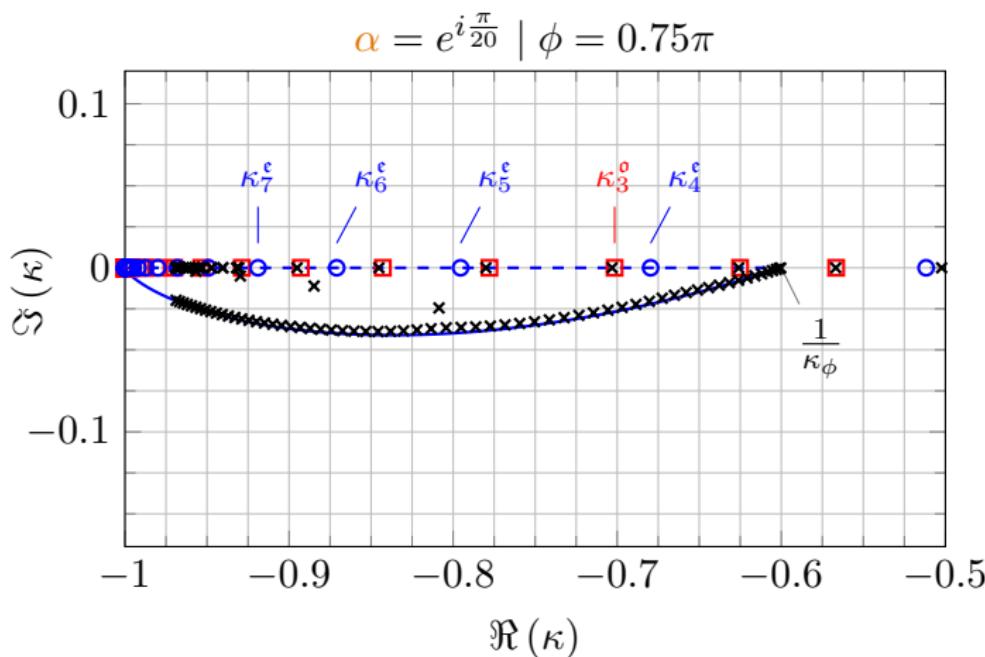


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(\times): computed ($N_h = 26345$, $R_\alpha = R/2$, $R_{\text{TR}} = 10^{-50} \cdot R$),
(\circ): unperturbed even eigenvalues κ_n^e , (\square): κ_n^o ,
(—): even critical curve I_c^e, α , (—): I_c^o, α .

Results: spectrum

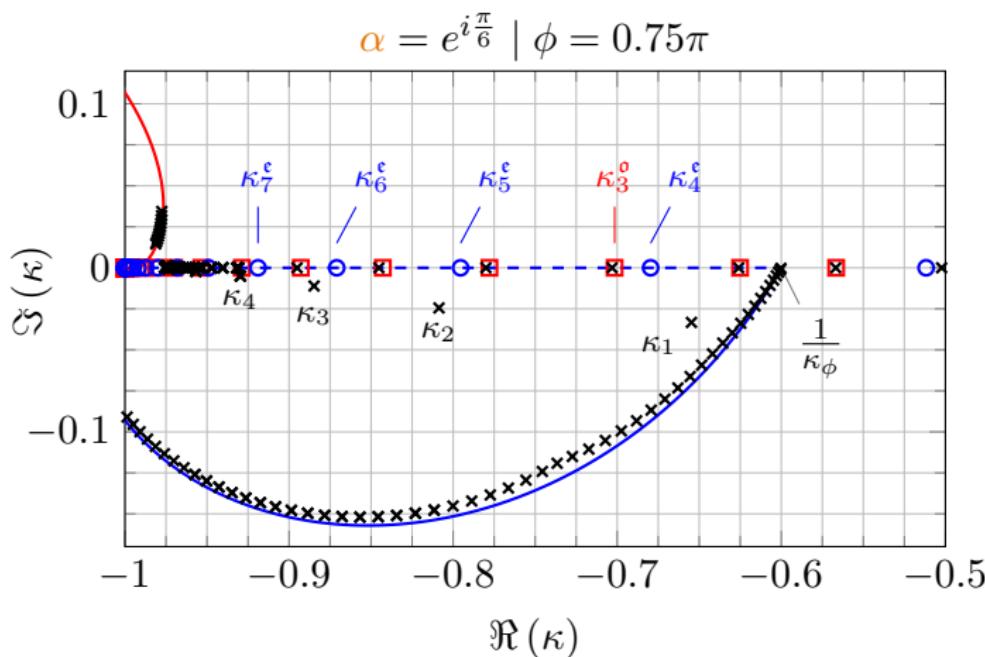


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Results: eigenfunctions

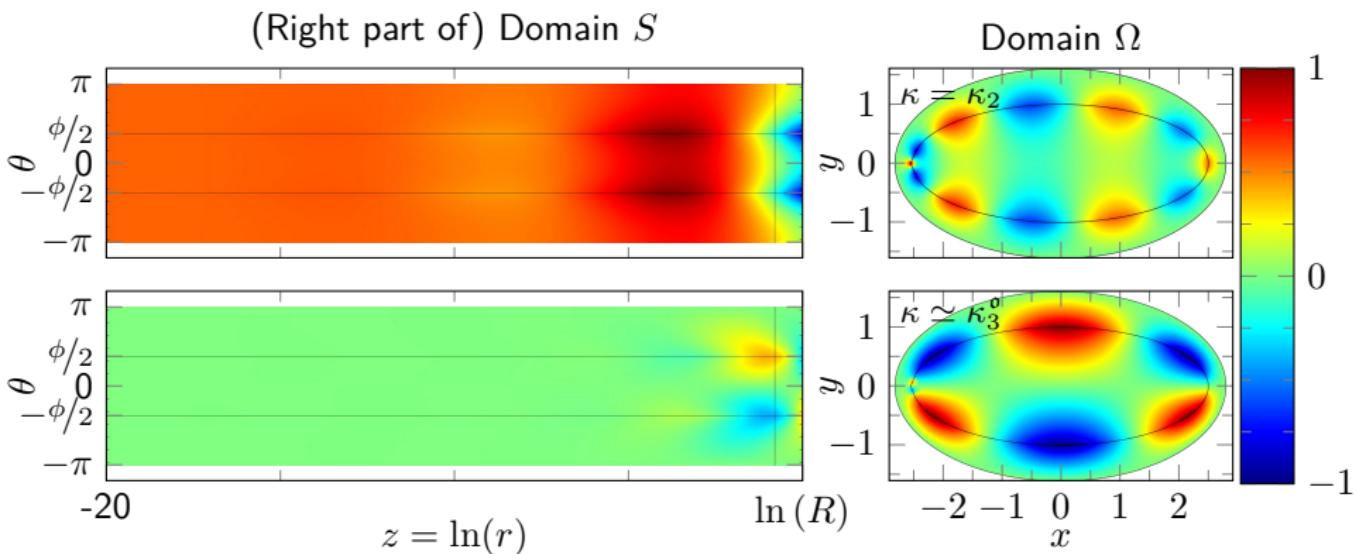


Fig. Eigenfunctions $\Re(u_\alpha)/\|u_\alpha\|_\infty$ of PEP- α with $\alpha = e^{i\frac{\pi}{6}}$.

(Top row) $\kappa = \kappa_2 \simeq 0.8086 - 0.02445i$, complex plasmonic resonance

(Bottom row) $\kappa \simeq 0.70313 - 8.0357 \cdot 10^{-8}i \simeq \kappa_3^o$, embedded eigenvalue.

Results: influence of corner angle

$$\alpha = e^{i\frac{\pi}{5}} \mid \phi = 0.86\pi$$

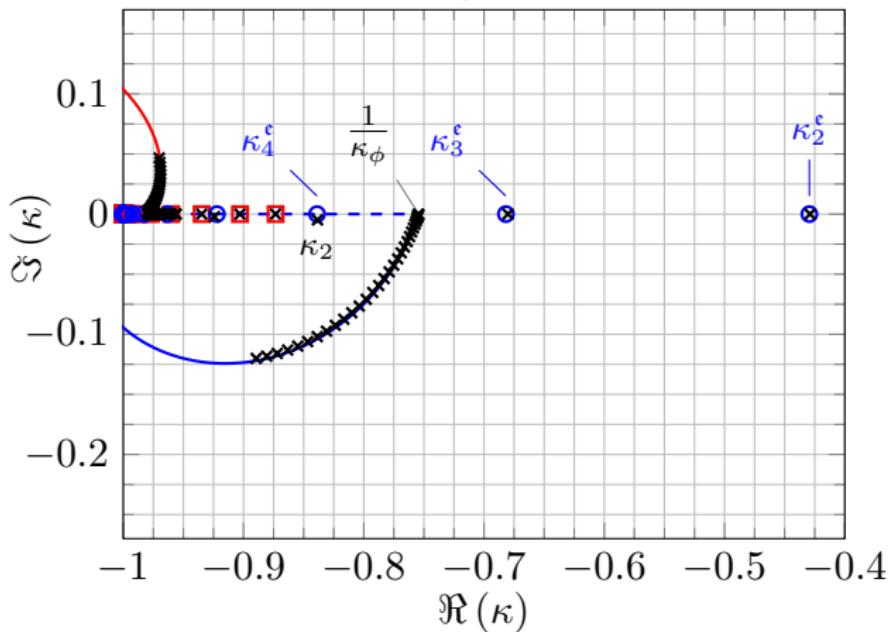


Fig. Spectrum for decreasing values of ϕ (increasing R).

(\times): computed ($N_h = 35695$, $R_\alpha = R/2$, $R_{\text{TR}} = 10^{-50} \cdot R$),
(\circ): unperturbed even eigenvalues κ_n^{ϵ} , (\square): κ_n^{σ} ,
(—): even critical curve $I_c^{\epsilon,\alpha}$, (—): $I_c^{\sigma,\alpha}$.

Results: influence of corner angle

$$\alpha = e^{i\frac{\pi}{5}} \mid \phi = 0.78\pi$$

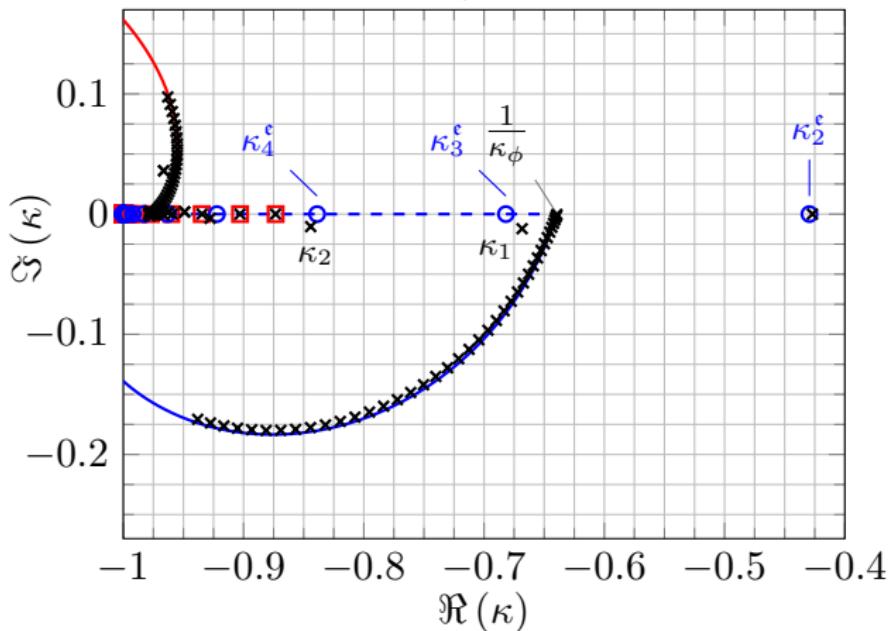


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(\times): computed ($N_h = 35695$, $R_\alpha = R/2$, $R_{\text{TR}} = 10^{-50} \cdot R$),
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(—): even critical curve $I_c^{\epsilon, \alpha}$, (—): $I_c^{\text{o}, \alpha}$.

Results: influence of corner angle

$$\alpha = e^{i\frac{\pi}{5}} \mid \phi = 0.63\pi$$

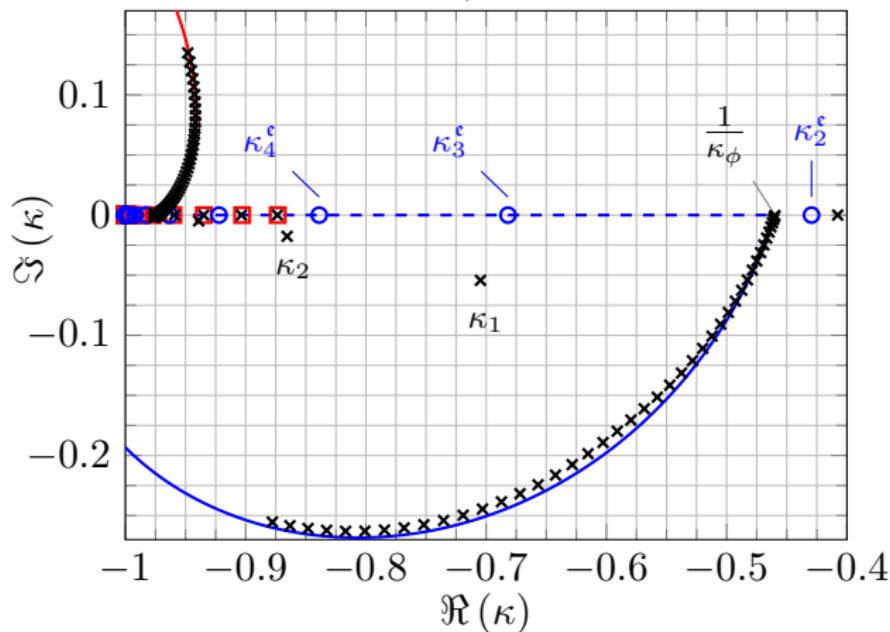


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Conclusions & outlook

Takeaways

- Complex plasmonic (CP) resonances
 - Defined using multivaluation of set of stable zeros $\kappa \mapsto \widehat{H}_\phi(\kappa)$
 - Analogous to complex resonances in scattering: "Infinity \Leftrightarrow Corner"
- Corner complex scaling (Bonnet-Ben Dhia, Carvalho, Chesnel, and Ciarlet 2016)
 - Yields a linear eigenvalue problem in κ , valid in uncovered region U_ϕ^α
- Numerical results
 - Meshing strategy for ellipse perturbed by corner
 - Corroborate mechanism described in (Li and Shipman 2019)

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Outlook

- Drop quasi-static assumption and solve the nonlinear problem in ω
- Interest of working with $\alpha(\kappa)$? (Nannen and Wess 2018)
- Extension to axi-symmetric particles $\Omega_m \subset \mathbb{R}^3$ (Helsing and Perfekt 2018)
- Computation of CP resonances with Boundary Element method?
(Helsing, Kang, and Lim 2017) (Helsing and Karlsson 2018)

[▶ Main TOC](#)[▶ Additional slides TOC](#)

Complex-scaling method for the plasmonic resonances of particles with corners

- 1 Introduction
- 2 Definition of complex plasmonic resonances
- 3 Definition of corner complex scaling
- 4 Strategy and numerical results
- 5 Conclusion

Thanks for your attention. Any questions?

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