

Recovery of remanent magnetizations using geometric measure theory, l_1 and total variation regularization

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The main problem we are interested in is the one of recovering the remanent magnetization \mathbf{M} of a rock sample from measurements of its magnetic field \mathbf{B} .

Using the magnetic potential Φ which satisfies

$$\begin{aligned}\mathbf{B} &= \mu_0 (\mathbf{M} - \text{grad } \Phi) \\ \Delta \Phi &= \text{div } \mathbf{M}\end{aligned}$$

we study the inverse problem for the recovery of \mathbf{M} when this magnetization is modeled by a vector valued measure.

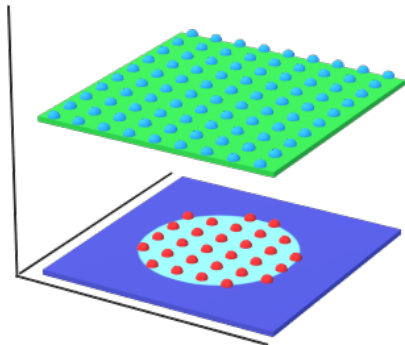
We use methods for recovering \mathbf{M} based on total variation regularization.

This work is part a collaborative project from the following:

- Eduardo Andrade Lima and Benjamin Weiss of the MIT Department of Earth, Atmospheric, and Planetary Sciences.
- Laurent Baratchart, Sylvain Chevillard and Juliette Leblond of the FACTAS, (formerly APICS) team in INRIA Sophia Antipolis.
- Doug Hardin and Edward Saff of Vanderbilt University

The measurements are usually obtained form a Scanning Magnetic Microscope (SMM) such as the instrument used by the research team in the MIT.

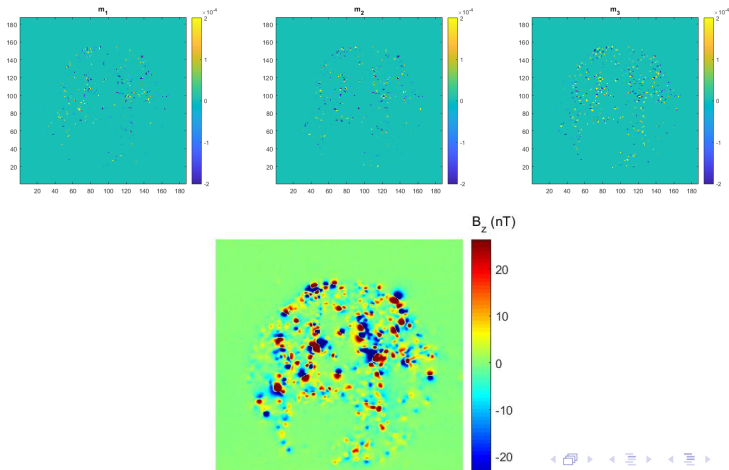
SMM setup



Q will denote the set on which we take the measurement.

S will denote a super set of the magnetization support.

Example of a reconstruction





The *magnetic moment* is a vector \mathbf{m} that measures the strength and direction of objects that generate magnetic fields.

A *magnetic dipole* can be viewed as the limit of a current loop that shrinks to a point with constant moment.

For a dipole at the point $\mathbf{y} \in \mathbb{R}^3$ with magnetic moment \mathbf{m} , the field at the point \mathbf{x} is

$$\mathbf{B}(\mathbf{x}) = -\frac{\mu_0}{4\pi} \left(\frac{\mathbf{m}}{|\mathbf{x} - \mathbf{y}|^3} - 3(\mathbf{x} - \mathbf{y}) \frac{(\mathbf{x} - \mathbf{y}) \cdot \mathbf{m}}{|\mathbf{x} - \mathbf{y}|^5} \right).$$

The *magnetization* of an object is a density for the magnetic moment.

We will model magnetizations by vector valued Borel measures.

$\mathcal{M}(S)$ denotes the space of Borel measures supported on $S \subset \mathbb{R}^3$.

Total variation of Magnetizations

Each $\mathbf{M} \in \mathcal{M}(S)^3$ is of the form $d\mathbf{M} = \mathbf{u}_{\mathbf{M}}d|\mathbf{M}|$, where $|\mathbf{M}|$ is a positive Borel measure supported on S and $|\mathbf{u}_{\mathbf{M}}| = 1$ $|\mathbf{M}|$ -a.e.

In this case the *Total Variation Norm* of a measure \mathbf{M} is

$$\|\mathbf{M}\|_{TV} := |\mathbf{M}|(\mathbb{R}^3).$$

Note that in the case where $\mathbf{M} = \text{grad } u$ then $\|\mathbf{M}\|_{TV}$ coincides with the total variation of u .

Definition Φ and \mathbf{B} as distributions on \mathbb{R}^3

For a magnetization \mathbf{M} we will define $\Phi(\mathbf{M}) \in L^1_{loc}(\mathbb{R}^3)$, the **scalar magnetic potential of \mathbf{M}** as the unique distribution that satisfies

$$\Delta\Phi = \operatorname{div} \mathbf{M}$$

and for any point \mathbf{x} not in the support of \mathbf{M}

$$\Phi(\mathbf{M})(\mathbf{x}) = \int (\operatorname{grad} \Gamma)(\mathbf{x} - \mathbf{y}) \cdot d\mathbf{M}(\mathbf{y}) = \frac{1}{4\pi} \int \frac{\mathbf{x} - \mathbf{y}}{|\mathbf{x} - \mathbf{y}|^3} \cdot d\mathbf{M}(\mathbf{y}),$$

where $\Gamma(\mathbf{x}) := -1/(4\pi|\mathbf{x}|)$ is the Newtonian kernel. Then we define the magnetic field $\mathbf{B}(\mathbf{M})$ generated by \mathbf{M} as

$$\mathbf{B}(\mathbf{M}) := \mu_0 (\mathbf{M} - \operatorname{grad} \Phi(\mathbf{M})).$$

Definition Forward operator

Consider a finite, positive Borel measure ρ with support contained in Q and let $\mathbf{A} : \mathcal{M}(S)^3 \rightarrow L^2(Q, \rho)$ be the operator defined by

$$\mathbf{A}(\mathbf{M})(\mathbf{x}) := \mathbf{v} \cdot \mathbf{B}(\mathbf{M})(\mathbf{x}), \quad \mathbf{x} \in Q. \quad (1)$$

Since $\mathbf{K}_{\mathbf{v}}$ is continuous on $\mathbb{R}^3 \setminus \{0\}$ and Q and S are positively separated, it follows that $B_{\mathbf{v}}$ is continuous on Q and consequently \mathbf{A} does indeed map $\mathcal{M}(S)^3$ into $L^2(Q, \rho)$.

We say that magnetizations $\mathbf{M}, \mathbf{N} \in \mathcal{M}(S)^3$ are S -equivalent if $\mathbf{B}(\mathbf{M})$ and $\mathbf{B}(\mathbf{N})$ agree on $\mathbb{R}^3 \setminus S$ in which case we write

$$\mathbf{M} \stackrel{S}{\equiv} \mathbf{N}.$$

A magnetization \mathbf{M} is said to be S -silent (or silent in $\mathbb{R}^3 \setminus S$) if it is S -equivalent to the zero magnetization; i.e., if $\mathbf{B}(\mathbf{M})$ vanishes on $\mathbb{R}^3 \setminus S$.

To fix notation, for $\mathbf{M} \in \mathcal{M}(S)^3$, let

$$\mathfrak{M}(\mathbf{M}) := \inf\{\|\mathbf{N}\|_{TV} : \mathbf{N} \stackrel{S}{\equiv} \mathbf{M}\}. \quad (2)$$

Extremal Problem (EP-1)

Given $\mathbf{M}_0 \in \mathcal{M}(S)^3$, find $\mathbf{M} \stackrel{S}{\equiv} \mathbf{M}_0$ such that $\|\mathbf{M}\|_{TV} = \mathfrak{M}(\mathbf{M}_0)$.

The second extremal problem involves minimizing the following functional defined for $\mathbf{M} \in \mathcal{M}(S)^3$, $f \in L^2(Q)$, and $\lambda > 0$, by

$$\mathcal{F}_{f,\lambda}(\mathbf{M}) := \|f - \mathbf{A}\mathbf{M}\|_{L^2(Q)}^2 + \lambda \|\mathbf{M}\|_{TV}. \quad (3)$$

Extremal Problem (EP-2)

Given $f \in L^2(Q)$, find $\mathbf{M}_\lambda \in \mathcal{M}(S)^3$ such that

$$\mathcal{F}_{f,\lambda}(\mathbf{M}_\lambda) = \inf_{\mathbf{M} \in \mathcal{M}(S)^3} \mathcal{F}_{f,\lambda}(\mathbf{M}). \quad (4)$$

We will call a closed set $S \subset \mathbb{R}^3$ a **slender set** if $\mathcal{L}_3(S) = 0$ and each connected component C of $\mathbb{R}^3 \setminus S$ satisfies $\mathcal{L}_3(C) = \infty$.

Lemma

Let $S \in \mathbb{R}^3$ be closed and slender.

Then \mathbf{M} is S -silent if and only if $\operatorname{div} \mathbf{M} = 0$.

S -silent not slender

Let S be the closed unit Euclidean ball centered at the origin, $\mathbf{M} = \mathcal{L}_3 \lfloor S$, $\mathbf{M} = (4\pi/3)^{-1} \mathbf{M} \mathbf{e}_1$ and $\mathbf{N} := \delta_0 \mathbf{e}_1$ where δ_0 is the Dirac delta at zero.

Using the mean value theorem it follows that both have magnetic potential

$$\frac{1}{4\pi} \frac{x_1}{|\mathbf{x}|^3}$$

so $\mathbf{M} - \mathbf{N}$ is S -silent.

However, $\mathbf{M} - \mathbf{N}$ is not divergence free.

Sparse in the sense of being purely 1-unrectifiable

A set $B \subset \mathbb{R}^n$ is **purely 1-unrectifiable** if $\mathcal{H}_1(f(\mathbb{R}) \cap B) = 0$ for every Lipschitz map $f : \mathbb{R} \mapsto \mathbb{R}^n$.

Theorem

Suppose $S \subset \mathbb{R}^3$ is a closed, slender set. If $\mathbf{M} \in \mathcal{M}(S)^3$ has support that is purely 1-unrectifiable and $\mathbf{N} \in \mathcal{M}(S)^3$ is S -equivalent to \mathbf{M} , then $\|\mathbf{N}\|_{TV} > \|\mathbf{M}\|_{TV}$ unless $\mathbf{N} = \mathbf{M}$.

We say that \mathbf{M} is **uni-directional** if $\mathbf{u}_{\mathbf{M}}$ is constant a.e. with respect to $|\mathbf{M}|$.

Theorem

Let $S = \bigcup_{i=1}^n S_i$ for some disjoint closed sets S_1, S_2, \dots, S_n in \mathbb{R}^3 and suppose that S is either compact or slender. Let $\mathbf{M} \in \mathcal{M}(S)^3$ be such that $\mathbf{M}_i := \mathbf{M} \lfloor S_i$ is uni-directional for $i = 1, 2, \dots, n$.

If $\mathbf{N} \in \mathcal{M}(S)^3$ is S -equivalent to \mathbf{M} , then $\mathbf{N}_i := \mathbf{N} \lfloor S_i$ and \mathbf{M}_i are S_i -equivalent for $i = 1, 2, \dots, n$, moreover

$$\|\mathbf{M}\|_{TV} \leq \|\mathbf{N}\|_{TV}, \quad (5)$$

with equality in (5) if and only if \mathbf{N}_i is uni-directional in the same direction as \mathbf{M}_i for $i = 1, 2, \dots, n$. Furthermore, if S is slender and equality holds in (5), then $\mathbf{M} = \mathbf{N}$.

Oriented rectifiable curves

We call a Lipschitz mapping $\mathbf{c} : [0, \ell] \rightarrow \mathbb{R}^3$ a **rectifiable curve** and let $\mathbf{C} := \mathbf{c}([0, \ell])$ denote its image.

If \mathbf{c} is an arclength parametrization of \mathbf{C} ; i.e., if \mathbf{c} satisfies

$$\mathcal{H}_1(\mathbf{c}([\alpha, \beta])) = \beta - \alpha, \quad \forall [\alpha, \beta] \subset [0, \ell],$$

then we call \mathbf{c} an **oriented rectifiable curve**.

We will call $\mathbf{R} \in \mathcal{M}(S)^3$ defined through the relation

$$\langle \mathbf{R}, \mathbf{f} \rangle = \int_0^\ell \mathbf{f}(\mathbf{c}(t)) \cdot \mathbf{c}'(t) dt,$$

a **curve magnetization**.

Loop silent magnetizations

An oriented rectifiable curve will be called a **rectifiable Jordan curve** if it is injective on $[0, l)$ and such that $c(0) = c(l)$.

And its respective curve magnetization \mathbf{M} is said to be a **loop silent magnetization**.



Any such magnetization is divergence free by the Gauss Green Theorem in the measure theoretical case.

Decomposition of divergence free magnetizations

Let $\mathcal{C} \subset \mathcal{M}(S)^3$ denote the collection of oriented rectifiable curves with topology inherited from $\mathcal{M}(S)^3$.

Suppose $\operatorname{div} \mathbf{M} = 0$ (as a distribution).

Smirnov [1, Theorem A] shows that \mathbf{M} can be decomposed into elements from \mathcal{C} .

In particular, it can be proven that there is a positive Borel measure ρ on \mathcal{C} such that

$$\mathbf{M}(B) = \int \mathbf{R}(B) \, d\rho(\mathbf{R}),$$

We will only consider planar magnetizations, that is magnetizations supported on a closed $S \subset \mathbb{R}^2 \times \{0\}$.

Recall that any such S is a slender set.

Theorem

Let S be a closed subset of \mathbb{R}^2 and $\mathbf{M} \in \mathcal{M}(S)^3$ a silent magnetization. Then, for a.e. $t \in \mathbb{R}$ there is a countable family $\{\mathbf{N}_m^t\}$ of silent loop magnetizations such that

$$\|\mathbf{M}\|_{TV} = \int_{\mathbb{R}} \left(\sum \|\mathbf{N}_m^t\|_{TV} \right) dt < +\infty, \quad (6)$$

and for any Borel set $B \subset \mathbb{R}^2$,

$$\mathbf{M}(B) = \int_{\mathbb{R}} \sum \mathbf{N}_m^t(B) dt. \quad (7)$$



A set is said to be **tree-like** if it contains no rectifiable Jordan curves.

Corollary

Let S be a closed subset of \mathbb{R}^2 and $\mathbf{M} \in \mathcal{M}(S)^3$ be S -silent. Then, the support of \mathbf{M} contains a rectifiable Jordan curve. Hence, if S is tree-like the only S -silent magnetization is the zero magnetization.

Regularized inverse problem

The problem we refer to EP2 has the following discrete analog:

$$\arg \min_{\mathbf{m}} \|\mathbf{b} - \mathbf{A}\mathbf{m}\|_2^2 + \lambda \sum_i \|\mathbf{m}_i\|_2.$$

$\|\mathbf{b} - \mathbf{A}\mathbf{m}\|_2^2$ measures the distance between the original field and the one generated by the reconstruction.

$\sum_i \|\mathbf{m}_i\|_2$ is used to enforce sparsity.

Illustration for purely 1-unrectifiable recovery

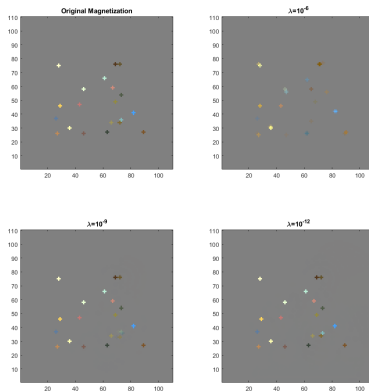
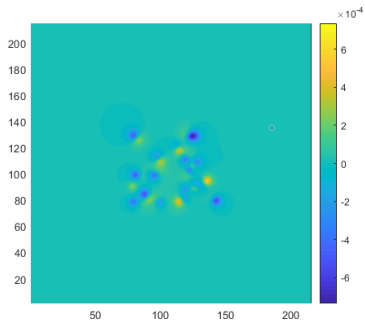
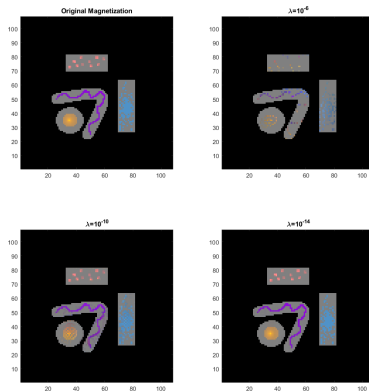
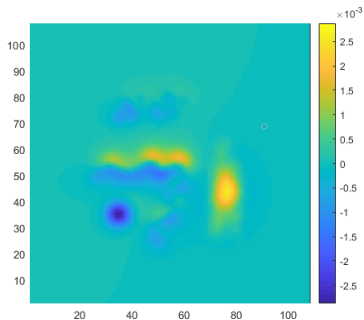
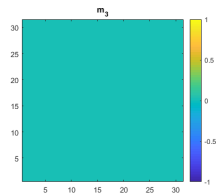
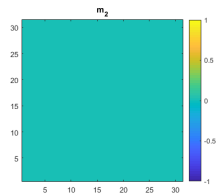
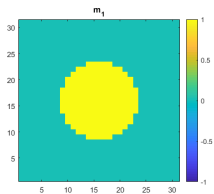


Illustration for Unidirectional recovery



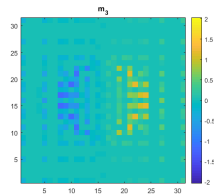
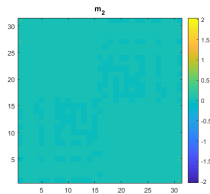
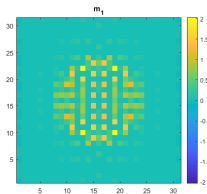
Unidirectional recovery in detail

Original



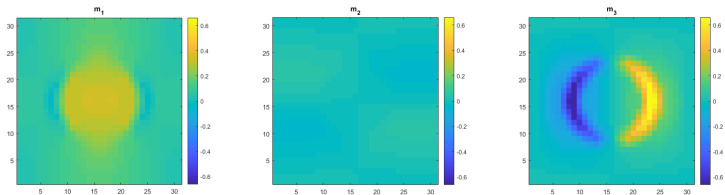
Unidirectional recovery in detail

l_1 with $6e-4$ relative residual



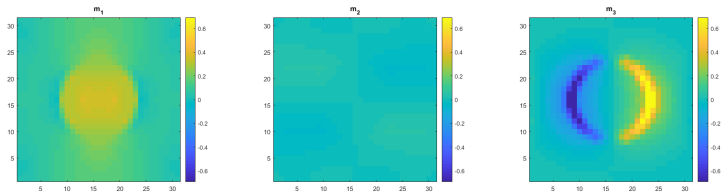
Unidirectional recovery in detail

TV with $4e-4$ relative residual (separated norm)



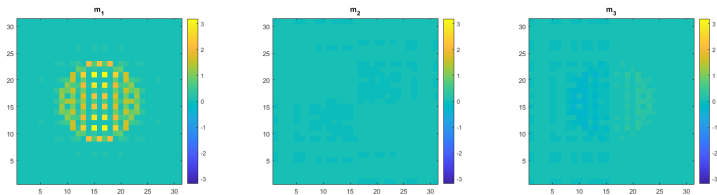
Unidirectional recovery in detail

TV with $4e-5$ relative residual (l_2 norm of components)



Unidirectional recovery in detail

Simulating full measurements, $5e-3$ residual.



Thank you



S. K. Smirnov.

Decomposition of solenoidal vector charges into elementary solenoids, and the structure of normal one-dimensional flows.

St. Petersburg Math. J., 5:841–867, 1994.



Pertti Mattila.

Geometry of sets and measures in Euclidean spaces, volume 44 of *Cambridge Studies in Advanced Mathematics*.

Cambridge University Press, Cambridge, 1995.

Fractals and rectifiability.